

Higher Representation Theory for Excitations in Gauge Theories

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Introduction

Today's message

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- Representation theory for **excitations**
 - **Higher groupoid representation**
 - Depends on **microscopic symmetry** and **infrared phase**
 - **Systematic** classification of excitations
 - **Higher** representation $\Rightarrow$ **extended** excitations
- Meaning of "gauge charge"
- ("Non-invertible" symmetries $\Rightarrow$ higher strip algebra representation)

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Representation theory in $1+1d$ (Cordova, Holfester, and Ohmori 2024)



Clay Córdova



Nicholas Holfester

Symmetry Preserved Phase

- G : finite group microscopic symmetry
- A single vacuum $|0\rangle$.
- Symmetry preserves the vacuum $g|0\rangle = |0\rangle$.
- Excitations are (and thus whole Hilbert space is) in a representation:

$$g \cdot |e\rangle = \rho_e(g)|e\rangle,$$

where $\rho_e(g)$ is a representation matrix.

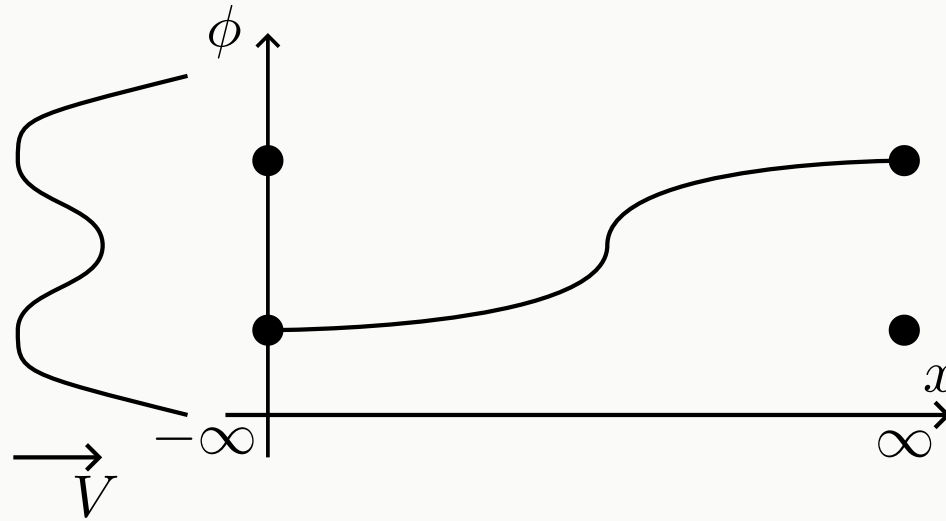
- Excitation $|e\rangle$ is classified by its **irreducible representation** of G .

Completely Broken Phase

- G : finite group (e.g. $\mathbb{Z}_2$, S_3 , $\dots$)
- Complete spontaneous symmetry breaking: $G \rightsquigarrow \{*\}$.
- $|G|$ degenerate vacua.
- $v \in G \Rightarrow$ vacuum $|v\rangle$.
 - Microscopic symmetry $g \in G$ acts as $g|v\rangle = |gv\rangle$.

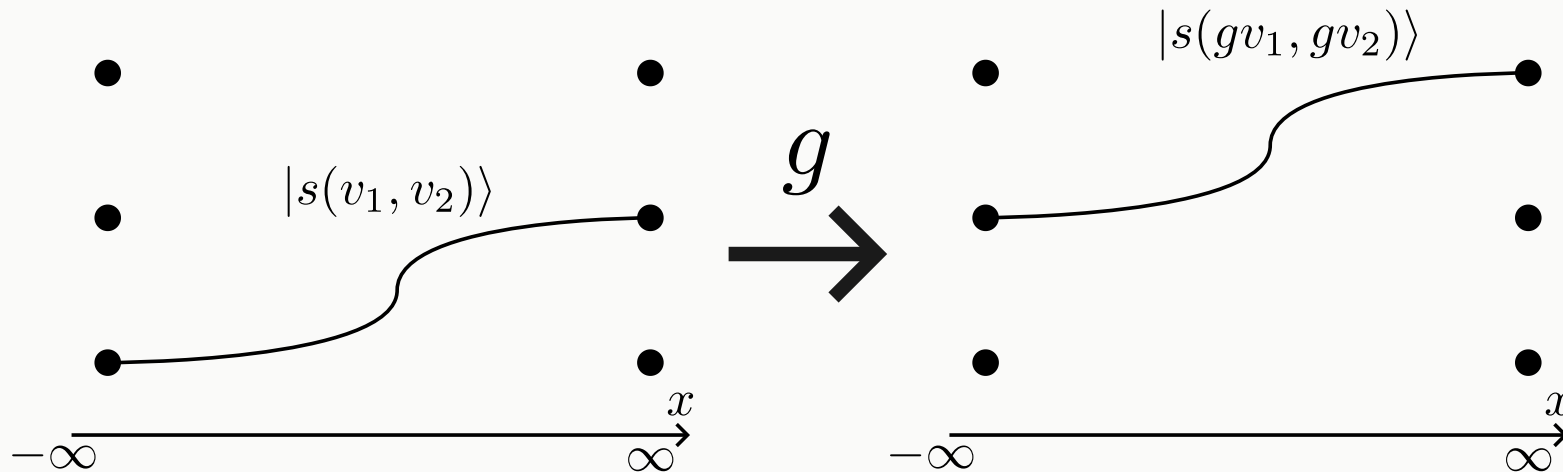
Solitons in SSB phase

- Soliton $|s(v_1, v_2)\rangle$ interpolates between v_1 and v_2 .
- $g|s(v_1, v_2)\rangle = |s(gv_1, gv_2)\rangle$.
- G -action generates **$|G|$ -fold multiplet**.



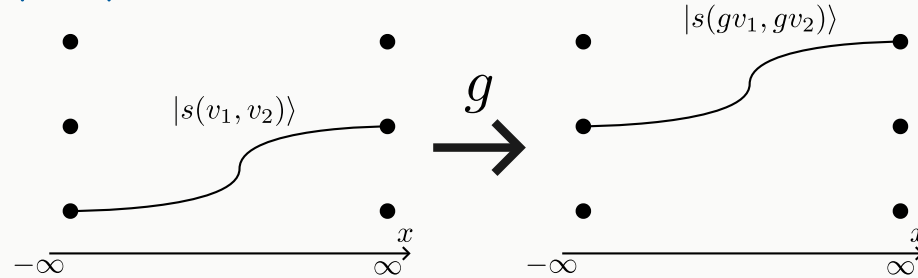
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Symmetry is "non-linearly represented".

A better understanding is that states in the phase are in (linear) representations of a **groupoid algebra**.

Groupoid Algebra

- The groupoid algebra $\mathbb{C}[X // G]$: a vector space spanned by $\overset{g}{\leftarrow} v$:
 - $v \in X$ is a vacua (at the left ∞ , say).
 - i.e. $\mathbb{C}[X // G] \ni \sum_{g,v} c_{g,v} [\overset{g}{\leftarrow} v]$
 - Multiplication: $[\overset{g_1}{\leftarrow} v_1] \cdot [\overset{g_2}{\leftarrow} v_2] = \delta_{v_1, g_2 v_2} [\overset{g_1 g_2}{\leftarrow} v_2]$.
- **Representation of groupoid algebra** (R, ρ) is an algebra hom $\rho : \mathbb{C}[X // G] \rightarrow \text{GL}(R)$ for a Hilbert space R .
- For $X = *$, (R, ρ) is a usual group representation.

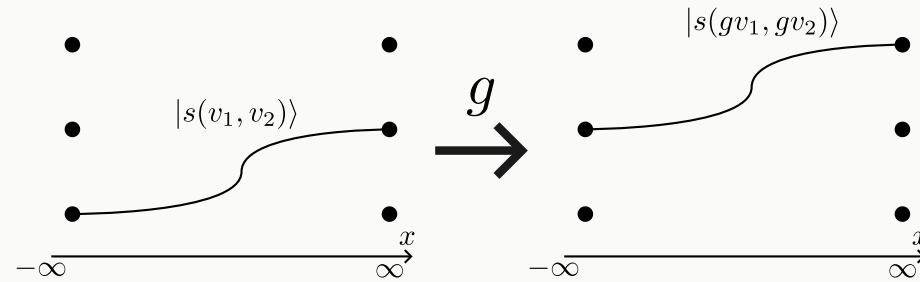
Groupoid for SSB

- For the SSB phase where G is completely broken, the relevant groupoid is $G^2 \parallel G$, where the action is $(v_1, v_2)g = (v_1g, g^{-1}v_2)$.
- $(v_1, v_2) \in G^2$: vacua at left and right infinity.
- Multiplication in $\mathbb{C}[G^2 \parallel G]$:

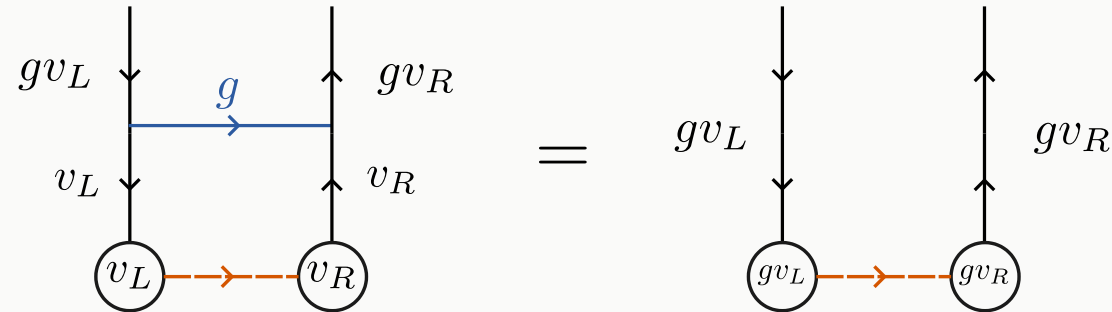
$$\begin{array}{c}
 \begin{array}{ccccc}
 \begin{array}{c} g_1 v_{L,1} \downarrow \\ v_{L,1} \downarrow \end{array} & \xrightarrow{g_1} & \begin{array}{c} \uparrow g_1 v_{R,1} \\ \uparrow v_{R,1} \end{array} & \bullet & \begin{array}{c} g_2 v_{L,2} \downarrow \\ v_{L,2} \downarrow \end{array} & \xrightarrow{g_2} & \begin{array}{c} \uparrow g_2 v_{R,2} \\ \uparrow v_{R,2} \end{array}
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 = \prod_{i=L,R} \delta_{v_{i,1}, g_2 v_{i,2}} & \begin{array}{ccccc}
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 \end{array}
 \end{array}$$

Groupoid Algebra Representation

- $$\left[\overset{g}{\leftarrow} (v_L, v_R) \right] |s(v'_L, v'_R)\rangle = \delta_{v_L, v'_L} \delta_{v_R, v'_R} |s(gv_L, gv_R)\rangle$$

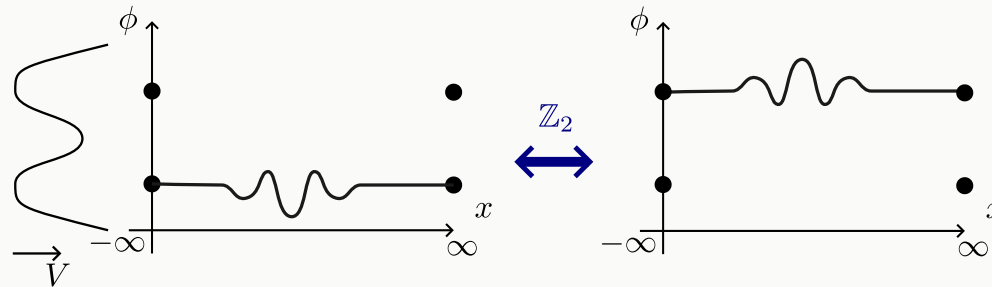


- Diagram:

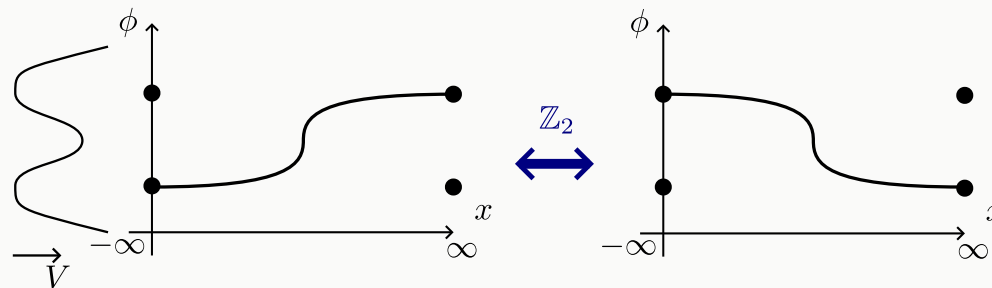


Groupoid Algebra Representation for $\mathbb{Z}_2$ SSB

- Irreps of $\mathbb{C}[\mathbb{Z}_2^2 // \mathbb{Z}_2]$: R_0, R_1 .
 - R_0 : "particle" sector : $(v_L, v_R) = (0, 0), (1, 1)$



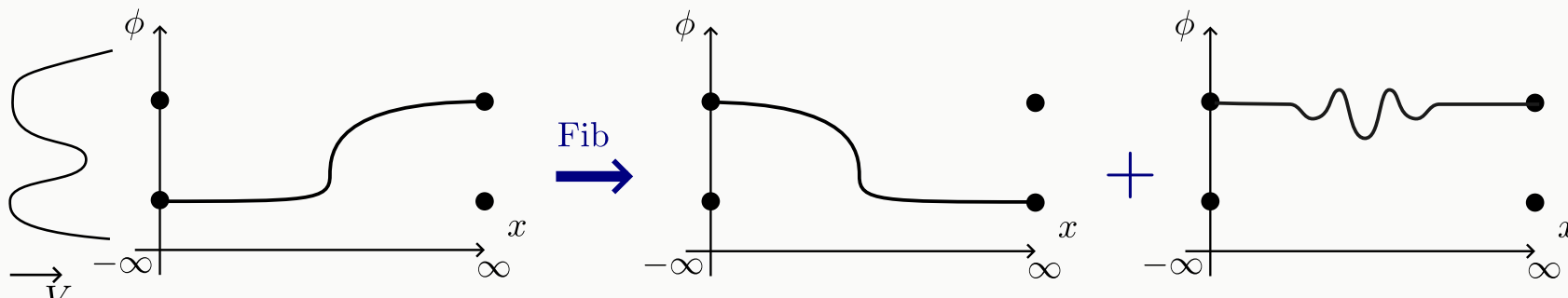
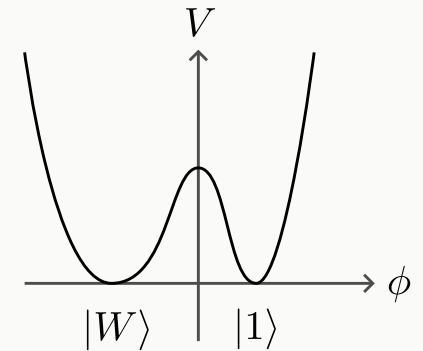
- R_1 : "soliton/antisoliton" sector : $(v_L, v_R) = (0, 1), (1, 0)$.



Rep. Theory for Solitons w/ Non-inv. Sym.

(Cordova, Holfester, and Ohmori 2024)

- Generalizes to **non-invertible symmetry**
- $\rightsquigarrow$ "**strip algebra**" (Kitaev and Kong 2012; Johnson-Freyd and Reutter WIP)
- E.g. Deformed tricritical Ising model (Zamolodchikov 1990)
 - SSB'en "Fibonacci" sym. (Chang et al. 2019)
 - a multiplet involving **both** particles and solitons!
(Córdova, García-Sepúlveda, and Holfester 2024)



General Theory

symbol	math	physics
$\mathcal{C}$	fusion category	UV symmetry/ Topological lines
$\mathcal{M}$	module category	IR gapped phase
$\mathbf{Str}_{\mathcal{C}}(\mathcal{M})$	weak Hopf alg.	Symmetry action on $\mathbb{R}^2$
$\mathbf{Rep}_{\mathcal{C}}(\mathcal{M}) \cong \mathbf{End}_{\mathbf{Mod}(\mathcal{C})}(\mathcal{M})$	Dual Category	Excitation multiplets

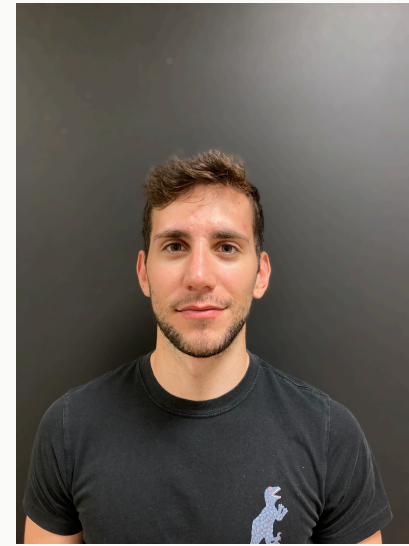
While $\mathbf{Str}_{\mathcal{C}}(\mathcal{M})$ is often cumbersome, one can often directly compute the representation category $\mathbf{Rep}_{\mathcal{M}}(\mathcal{C})$. E.g. full SSB: $\mathbf{Rep}_{\mathcal{C}}(\mathcal{C}) \cong \mathcal{C}$.

Representation Theory for Gauge Theories

Based on ([Gagliano, Grigoletto, and Ohmori 2025](#), WIP)



Finn Gagliano



Andrea Grigoletto

Higher Form Symmetry in Gauge theory

- pure YM with non-ab. gauge group $G \Rightarrow Z(G)^{(1)}$ **form symmetry**
 - e.g. $G_{\text{gauge}} = SU(N) : \mathbb{Z}_N^{(1)}$ one-form symmetry
 - $W_{\square}$ cannot be screened by gauge fields.

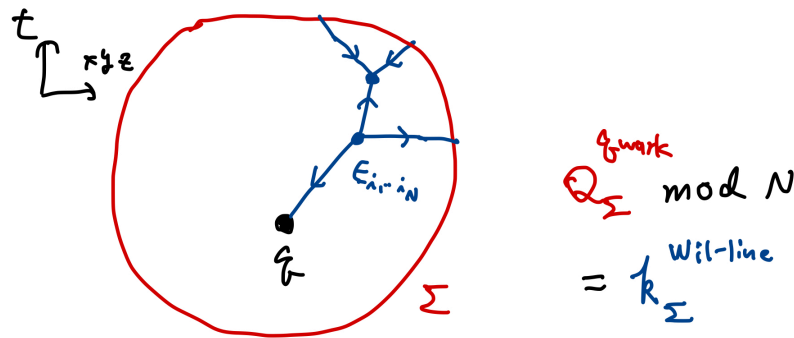
confinement (or not) $\Leftrightarrow \mathbb{Z}_N^{(1)}$ is preserved (or SSB'en).

Symmetry in single flavor QCD

- $G_{\text{gauge}} = \text{SU}(N) + \text{a (scalar) quark } q \text{ in } \square$
- $U(1)_q : q \mapsto e^{i\theta} q : \text{quark rotation}$
- $(\mathbb{Z}_N)_q = Z(G_{\text{gauge}}) \subset G_{\text{gauge}} : \text{gauged!}$
- $G_{\text{UV}} = U(1)_B = U(1)_q / \mathbb{Z}_N$: **Baryon number**
 - A local (gauge invariant) operator has **integer** Baryon charge.
 - $q_{i_1} \partial q_{i_2} \cdots \partial^N q_{i_N} \epsilon^{i_1, \dots, i_N}$

Symmetry Matching for Heavy Quark

- m_q : quark mass
 - $\mu \leq m_q: \mathbb{Z}_N^{(1)}$: **emergent** one-form symmetry
- Symmetry matching homomorphism: $U(1)_B \xrightarrow{\beta_N} \mathbb{Z}_N^{(1)}$
 - **Non-trivial!** (Bockstein map) c.f. (Seiberg and Seifnashri 2025)



Adjoint Higgs Phase

- $G_{UV} = U(1)_B \xrightarrow{\mu < m_q} \mathbb{Z}_N^{(1)} \xrightarrow{\mu < v_\Phi} *$
- "Exact sequence": $U(1)_q \rightarrow U(1)_B \rightarrow \mathbb{Z}_N^{(1)}$
 - " $\mathbb{Z}_N^{(1)} = U(1)_B / U(1)_q$ "
 - "break **up**" : $U(1)_B \rightsquigarrow U(1)_q$
- $\text{Rep}_{U(1)_B}(\text{Adj. Higgs}) \cong 3\text{Rep}(U(1)_q) \boxtimes_\omega 3\text{Vect}_{\mathbb{Z}_N^{(1)}}$
- **quark** : $\rho \in 3\text{Rep}(H = U(1)_q)$
- **center vortex** : $\sigma \in 3\text{Vect}(G/H = \mathbb{Z}_N^{(1)})$.

Summary

Message

Representation theory for excitations **depends on the phase**

- Discrete symmetry SSB phase $\Rightarrow$ groupoid algebra representation
- 1+1d SSB phase $\Rightarrow$ **solitons**
- (scalar) QCD
 - One-form symmetry pres / SSB
 $\Rightarrow$ **Representation theory** does not/ does contain quark

Prospects

Our work: **bosonic** theory in **gapped** and **zero-temperature** phases.

- Generalizations
 - Thermal Phase
 - Full QCD phase?
 - Gapless phases
 - Full gauge charge in Coulomb phase?
 - Quantum Gravity?

Prospects

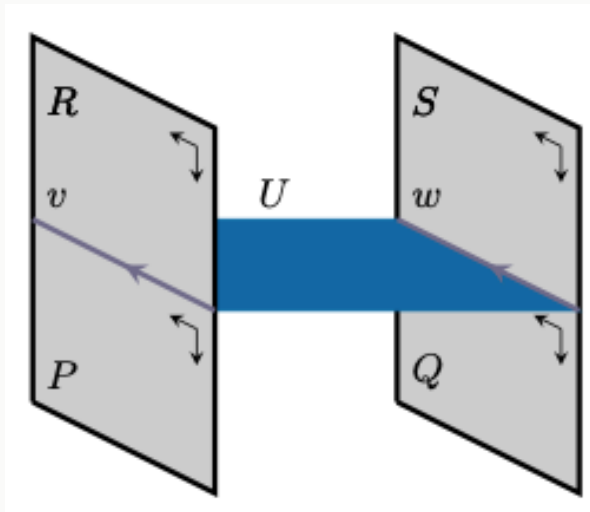
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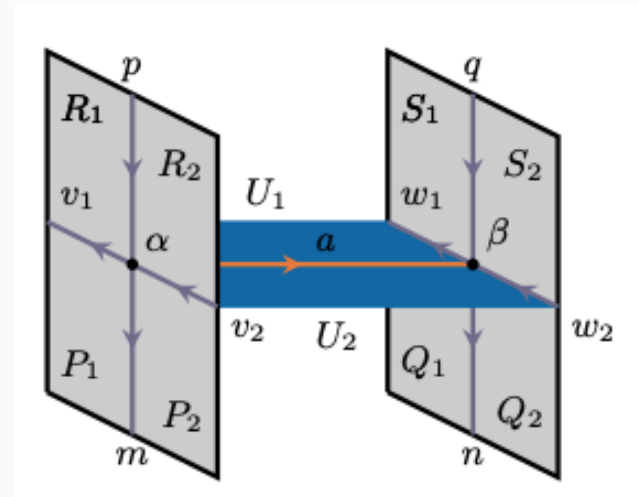
Thank you !!

Higher Strip Algebra

- In $d + 1$ QFT, higher strip algebra:
 - an weak Hopf algebra object in $d\mathbf{Vect}$.
 - $\text{Rep. theory} \cong \mathbf{Vect}_{G/H} \otimes_{\omega} \text{Rep}(H)$



object ($d = 2$)



1-morphism ($d = 2$)