



Quantum black hole mimicker: Redshift plays the role

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Image features of classical BHs:

- Bright ring
 - Highly lensed images on top of direct emission
 - Spherical photon orbits around the BH (photon sphere)



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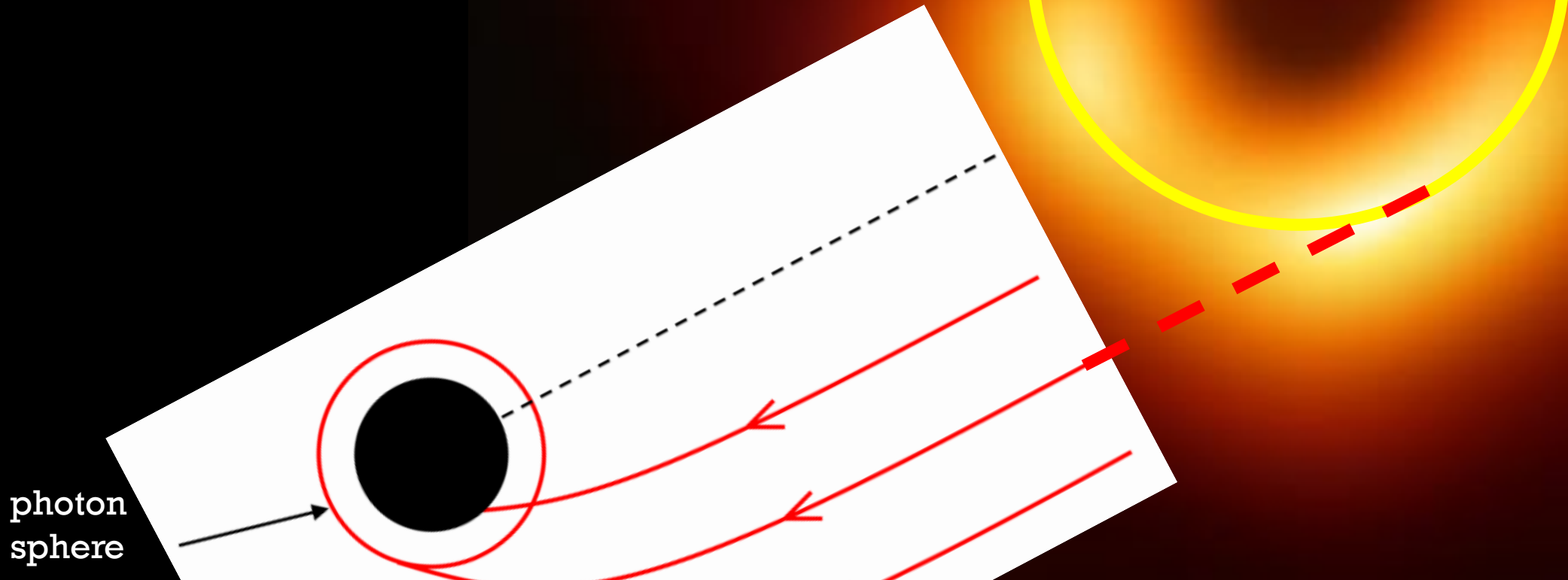


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 - Spherical photon orbits around the BH (photon sphere)
- Central brightness depression
 - Event horizon

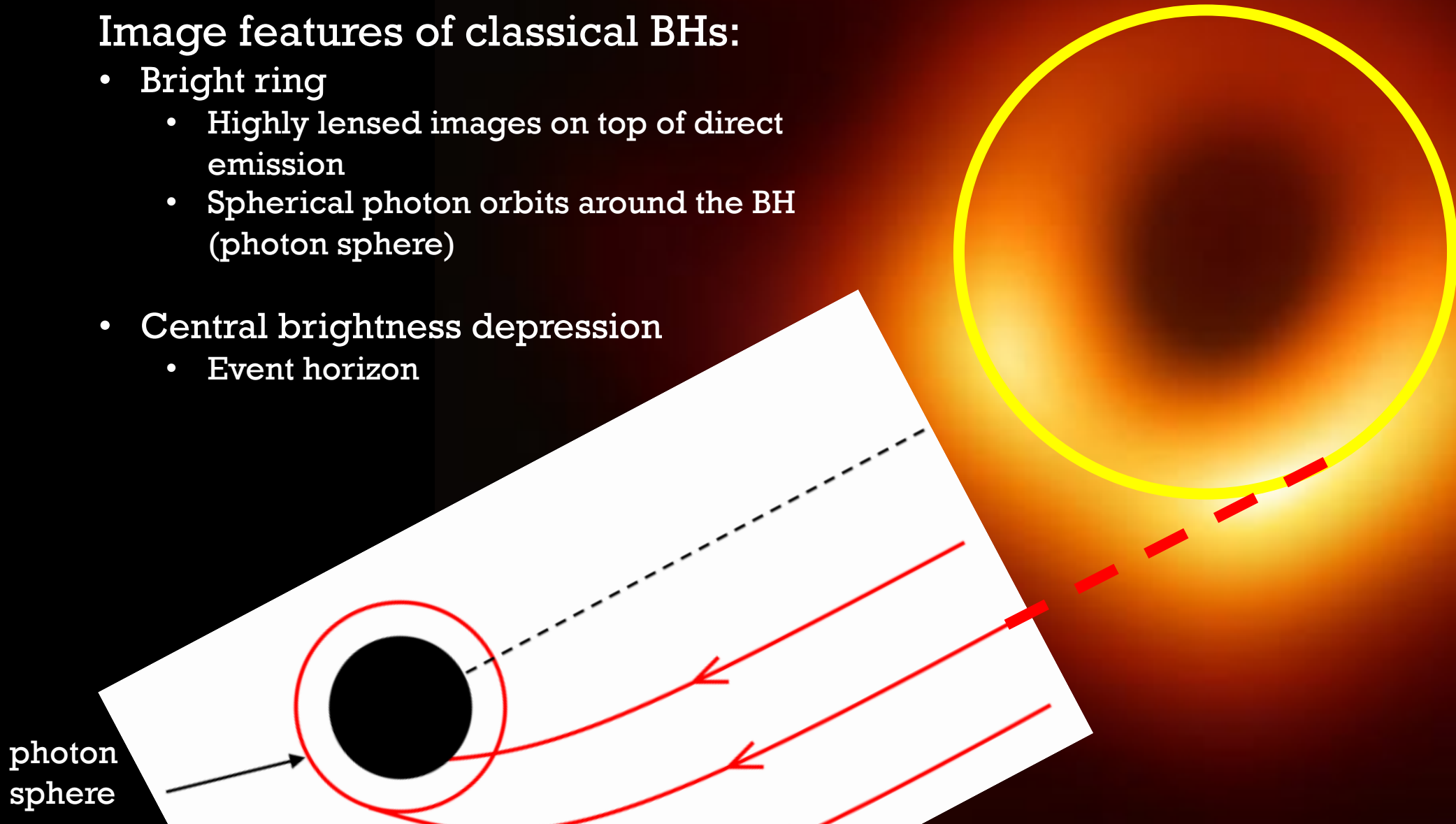
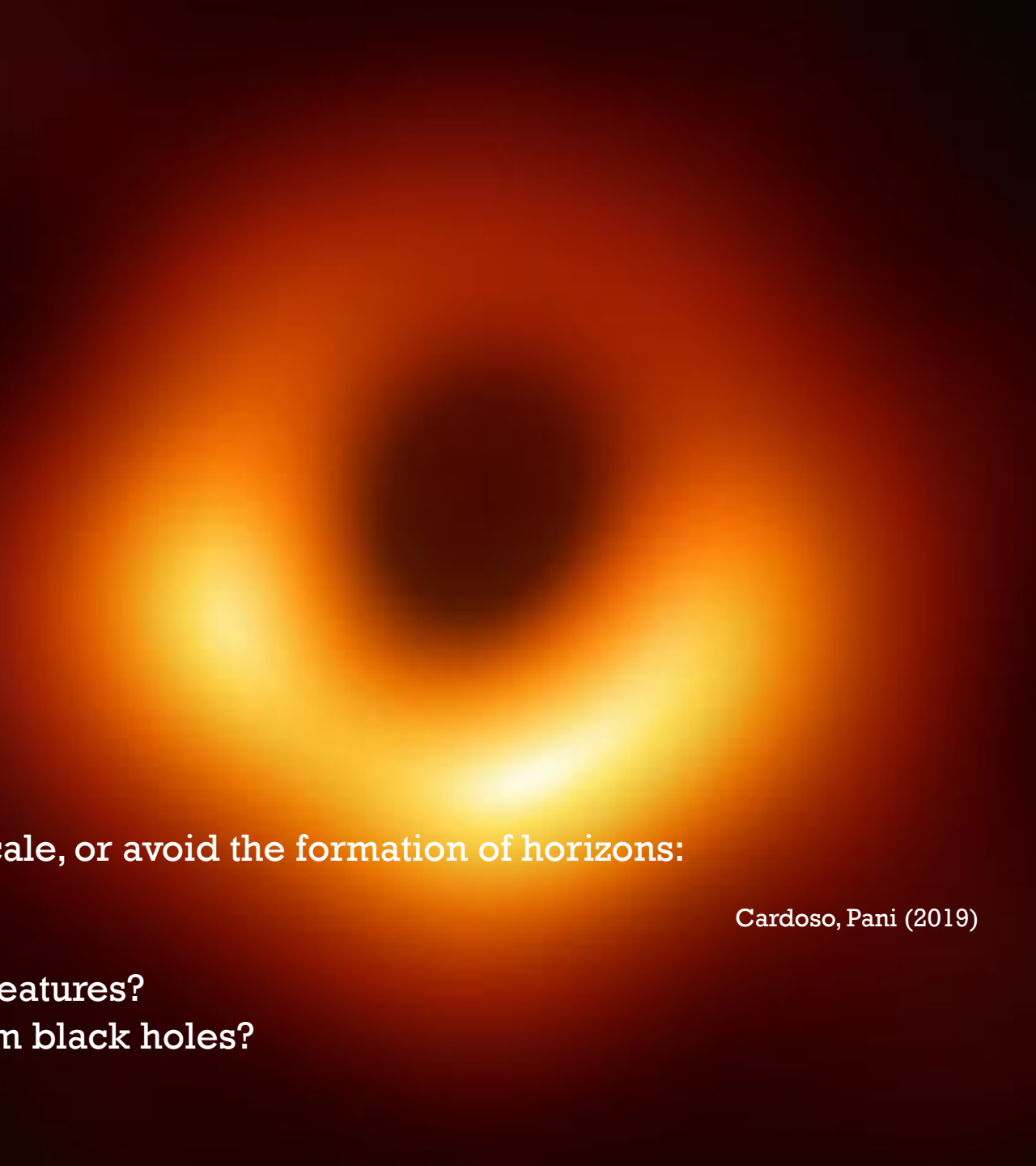


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Quantum effects may have horizon-scale, or avoid the formation of horizons:
Horizonless compact object (HCO)

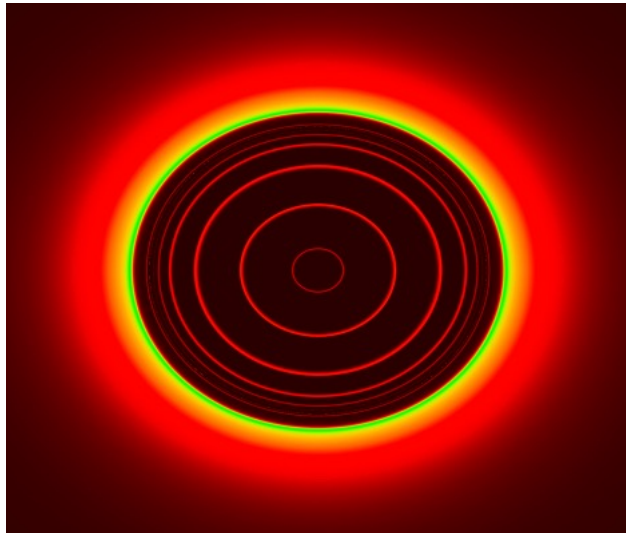
Cardoso, Pani (2019)

Can such objects cast similar image features?
If so, can it be a candidate for quantum black holes?

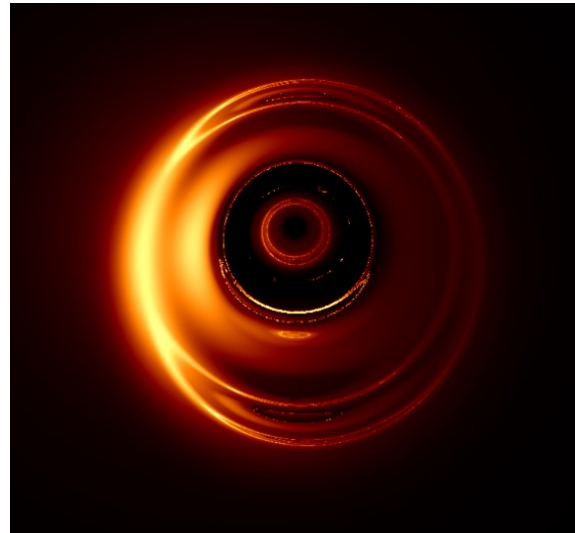
Image features of HCO

- Bright ring can appear if there is unstable photon sphere
- Additional inner rings
 - Reflection at the surface, internal propagation ...

Ohgami, Sakai (2015), Shaikh et al. (2019), Vincent et al. (2021), Wielgus et al. (2020)
Cunha et al. (2017), Rosa (2023) Mayerson, Vercnocke (2023), Shaikh et al. (2019)



Olmo et al. (2023)



Mayerson, Vercnocke (2023)

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- Two darkening effects may happen inside!!
 - Interaction between light rays and internal quantum structure
theory-dependent, challenging
 - Delayed by strong redshifts
depends only on metric, **crucial role in darkening images**



Quantum Horizonless Compact Object

- A HCO characterized by strong redshift inside

- Exterior (1):

$$g_{\mu\nu}^{(1)} = -\left(1 - \frac{a_0}{r}\right) dt^2 + \left(1 - \frac{a_0}{r}\right)^{-1} dr^2 + r^2 d\Omega^2, \quad r \geq R_{out} > a_0$$

- Interior dense region (2):

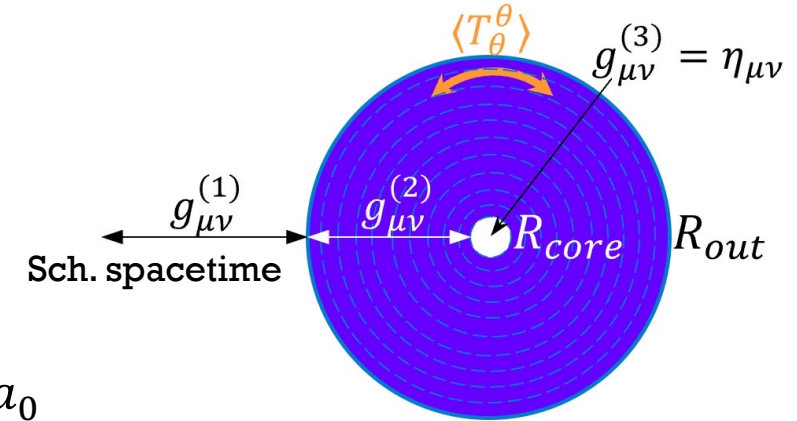
$$g_{\mu\nu}^{(2)} \approx -\frac{ka_0^2}{2r^2} e^{\frac{\eta(r^2 - R_{out}^2)}{2ka_0^2}} dt^2 + \frac{\eta^2 r^2}{2ka_0^2} dr^2 + r^2 d\Omega^2, \quad R_{core} < r < R_{out}$$

Kawai, Matsuo, Yokokura (2013), Kawai, Yokokura (2014)(2015)(2017)(2020)(2021), Yokokura (2022)(2023)

- Central core (3): $g_{\mu\nu}^{(3)} \approx \eta_{\mu\nu}, \quad r < R_{core}$

- A unified description of regions (2) and (3) has recently proposed

Livine, Yokokura (2025)



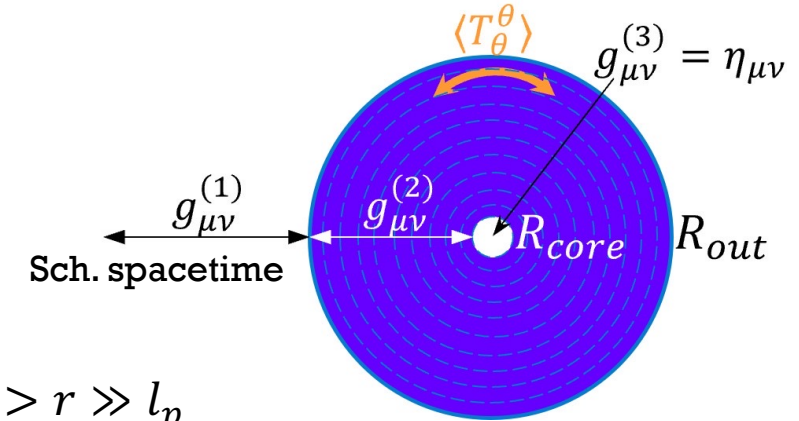
$$R_{out} \approx a_0 \left(1 + \frac{k}{2}\right)$$

Quantum Horizonless Compact Object

- Interior dense region:

$$g_{\mu\nu}^{(2)} \approx -\frac{ka_0^2}{2r^2} e^{\frac{\eta(r^2 - R_{out}^2)}{2ka_0^2}} dt^2 + \frac{\eta^2 r^2}{2ka_0^2} dr^2 + r^2 d\Omega^2$$

- Satisfy semi-classical Einstein eq. $G_{\mu\nu} = 8\pi G \langle \psi | T_{\mu\nu} | \psi \rangle$, where $R_{out} > r \gg l_p$
- Non-perturbative solution in \hbar
 - e.g. conformal matter field \rightarrow Trace anomaly $\rightarrow R \sim n l_p^2 R^2 \rightarrow$ Non-perturbative branch
 - $k \approx n l_p^2 / a_0^2$ with $n \gg 1$ (semi-classical)
 - $1 \leq \eta < 2$ effectively quantifies internal interactions
- Large tangential pressure
 - Highly anisotropic \rightarrow extreme compactness



Kawai, Yokokura (2020)

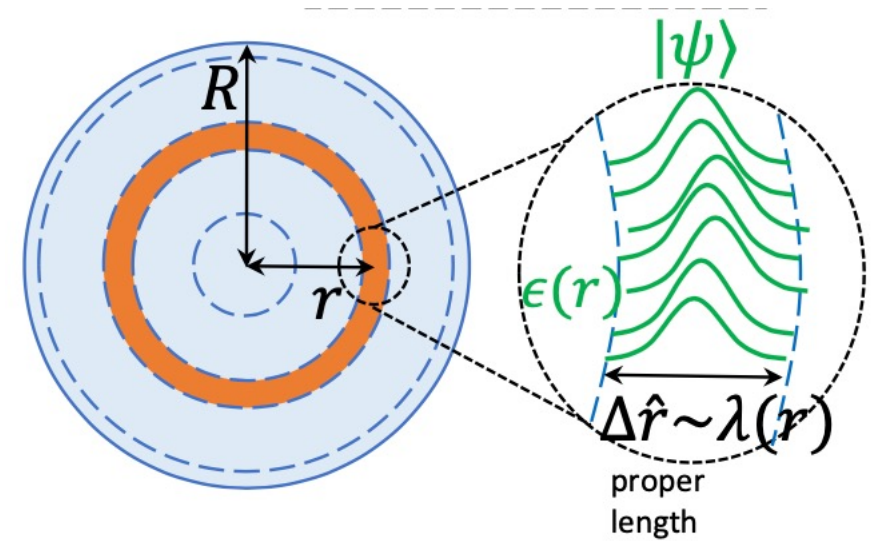
QHCO v.s. classical BH

- Interior dense region:

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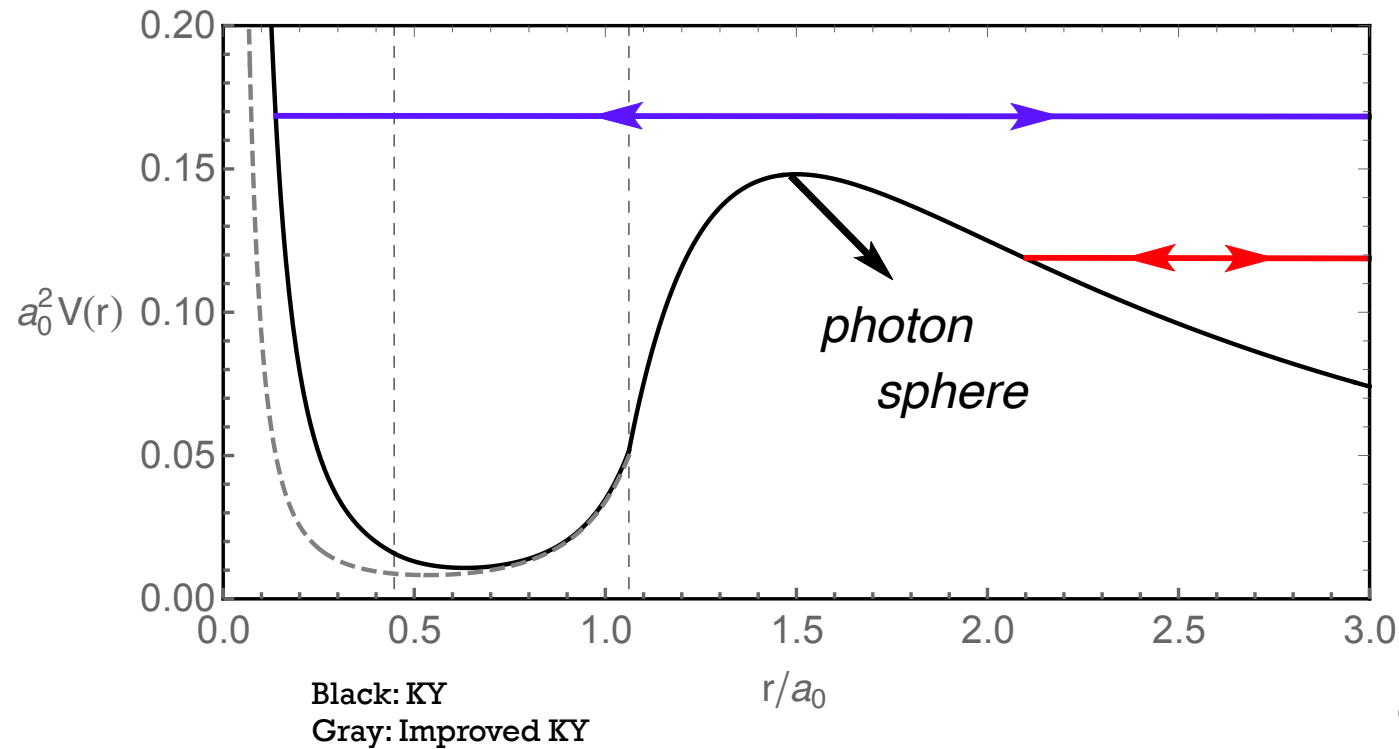
- Maximized entropy configuration
- Saturation conditions of entropy bounds
 - Entropy saturation: fix $g_{rr}(r)$
 - Positive pressure & $O(r^0)$ -curvature & Junction at R_{out} : fix $g_{tt}(r)$
- Exponentially strong redshift
- What does it look like?

Yokokura (2022)(2023)

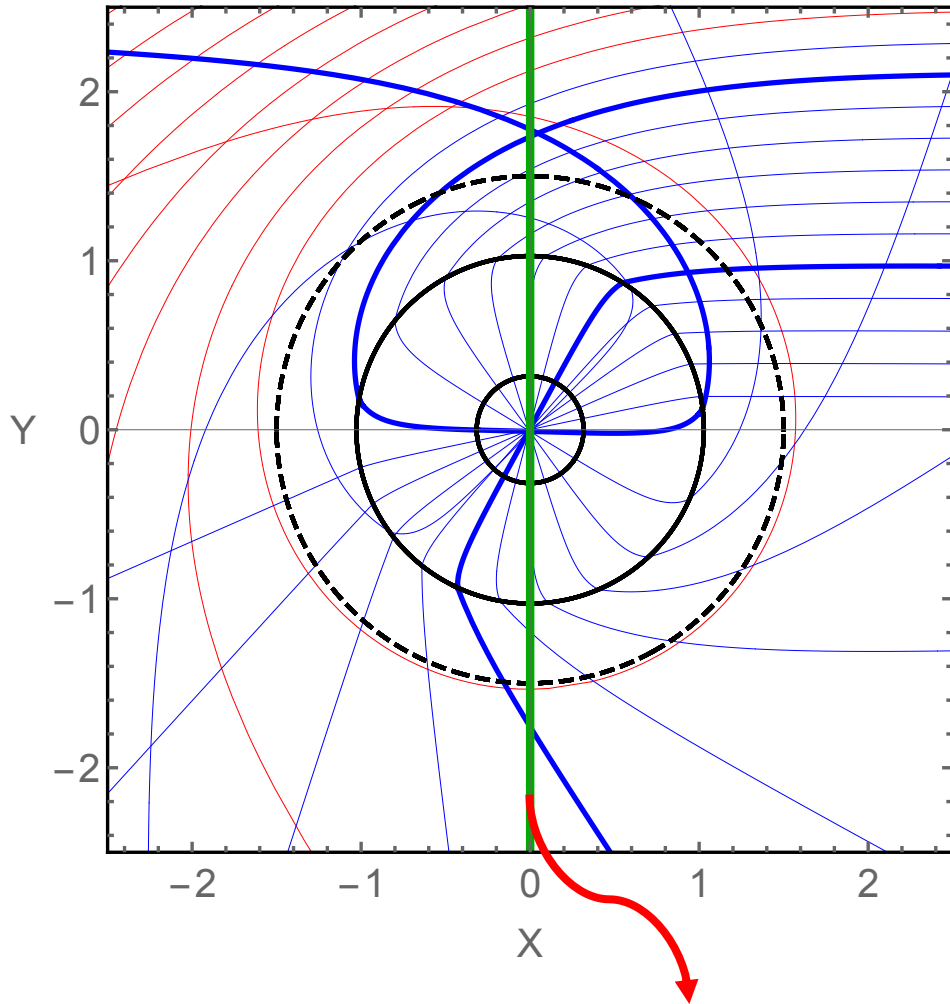


Light rays trajectories

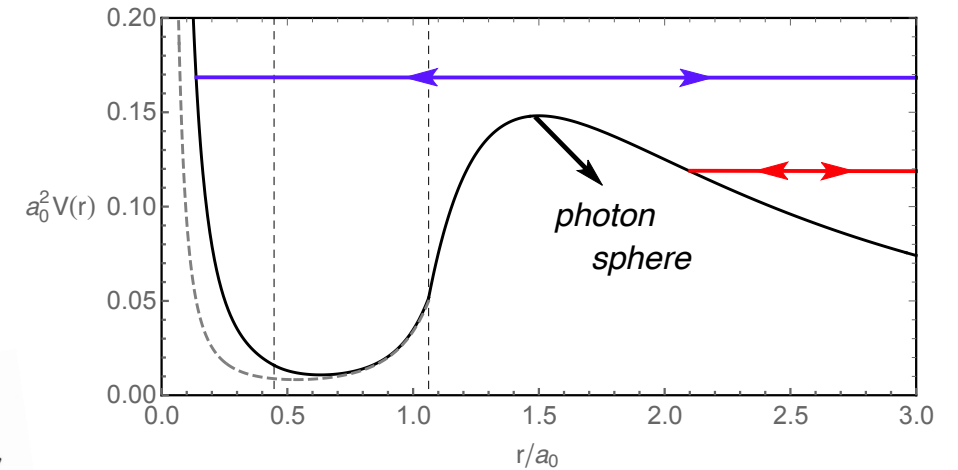
Effective potential: $\frac{C_i(r)}{L^2} \left(\frac{dr}{d\lambda} \right)^2 + V(r) = \frac{1}{b^2}$ $b = L/E$: impact parameter



Modeling QHCO images



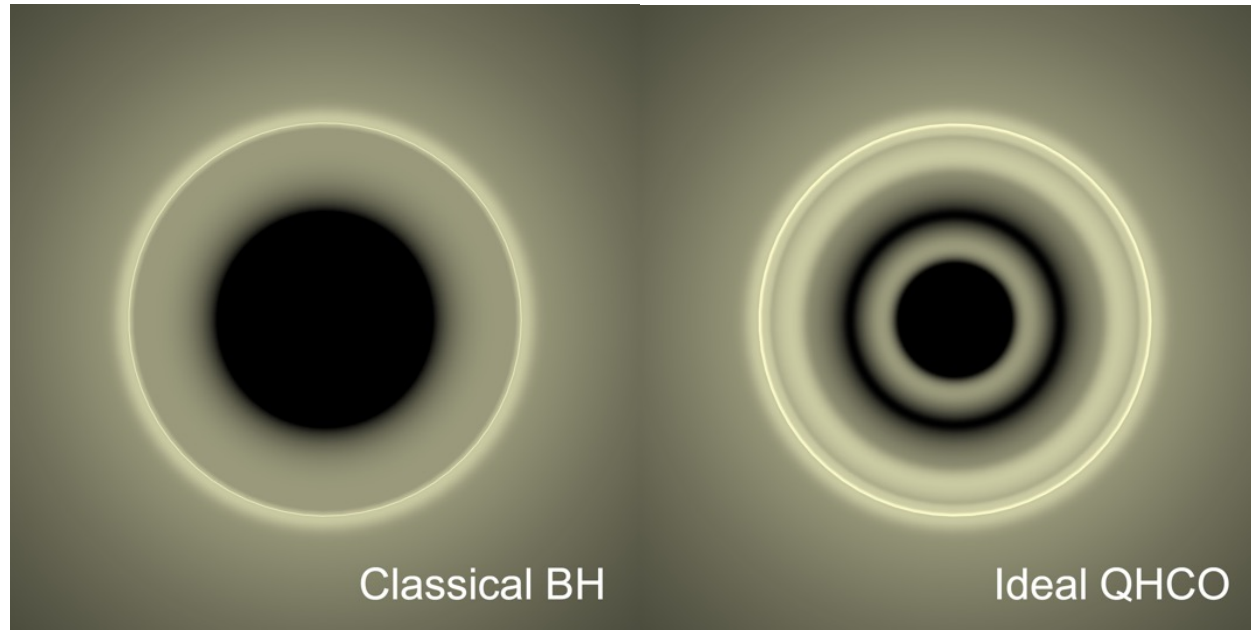
Geometrically thin, optically thin disk



1. Calculate trajectories, which are labeled by the impact parameter b
2. Identify the points where the trajectories cross the disk
3. Assume emission profile of the disk
4. For each trajectory, calculate the observed intensity $I_o(b)$

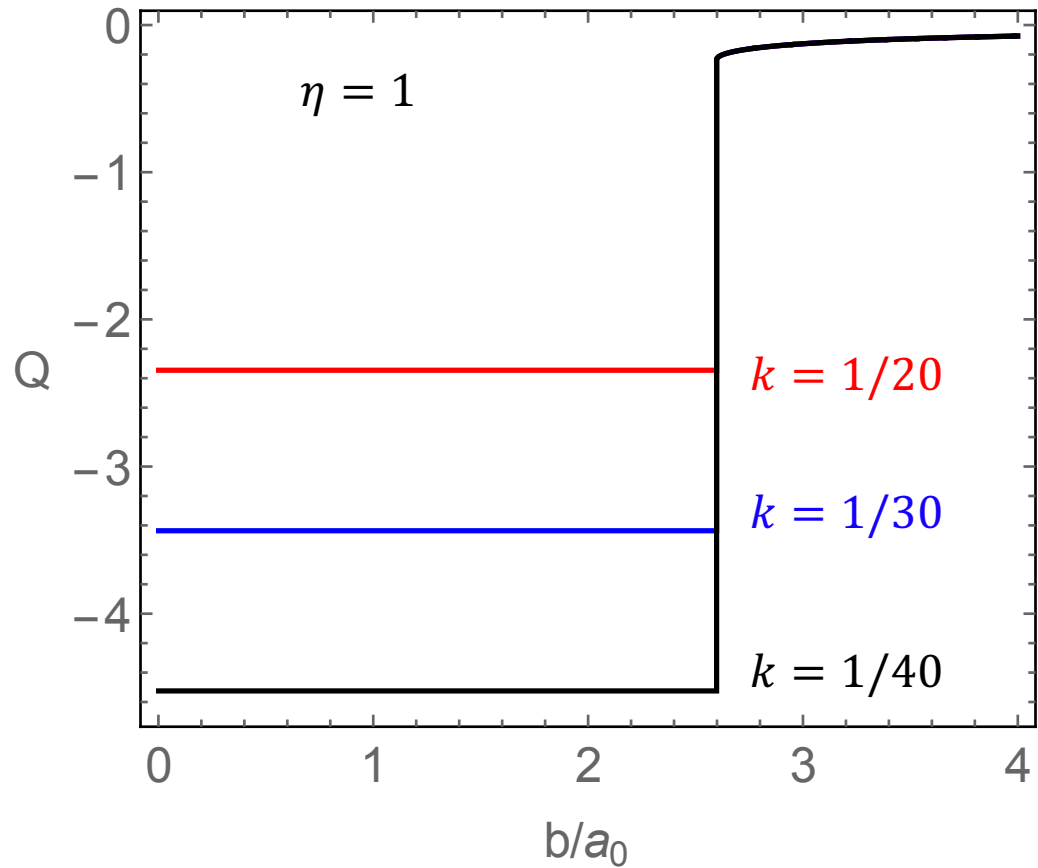
Ideal images

- Emission profile: emission down to the interior
- Assuming that the infalling light rays can be observed, **no matter how much time they may take with respect to the observer**



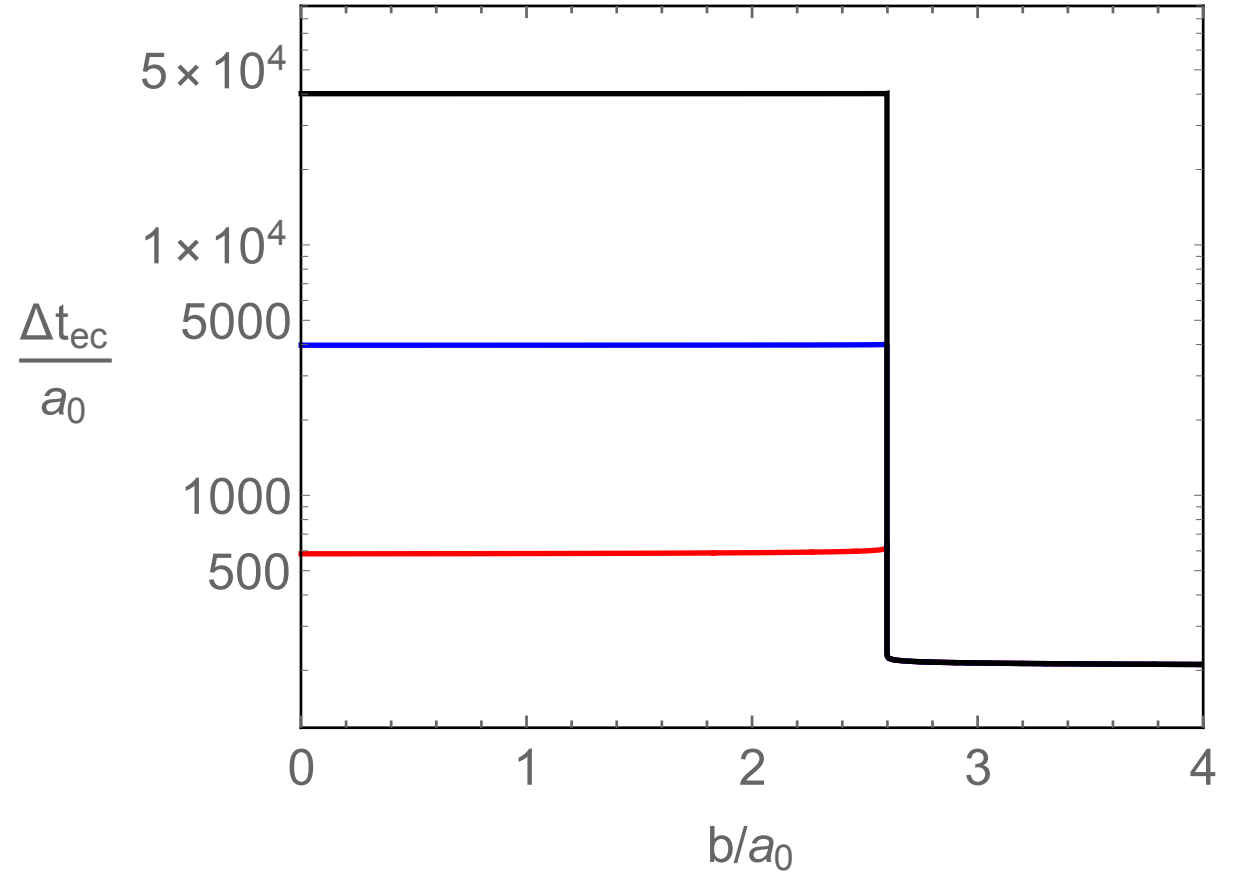
- Additional inner rings

Strong redshift and propagation time



$$Q \equiv \min[\log_{10} \sqrt{g_{tt}(\mathbf{x})/g_{tt}(\mathbf{x}_0)}]$$

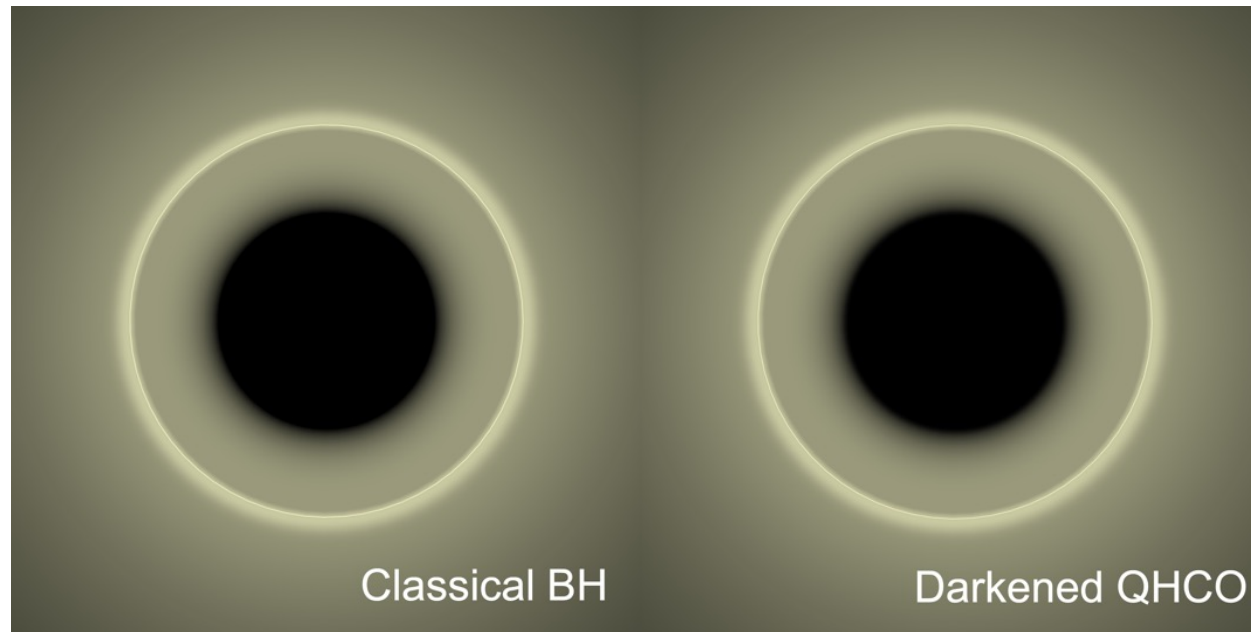
(maximized redshift inside)



$\Delta t_{ec} \sim a_0 e^{\eta/4k}$: elapsed coordinate time

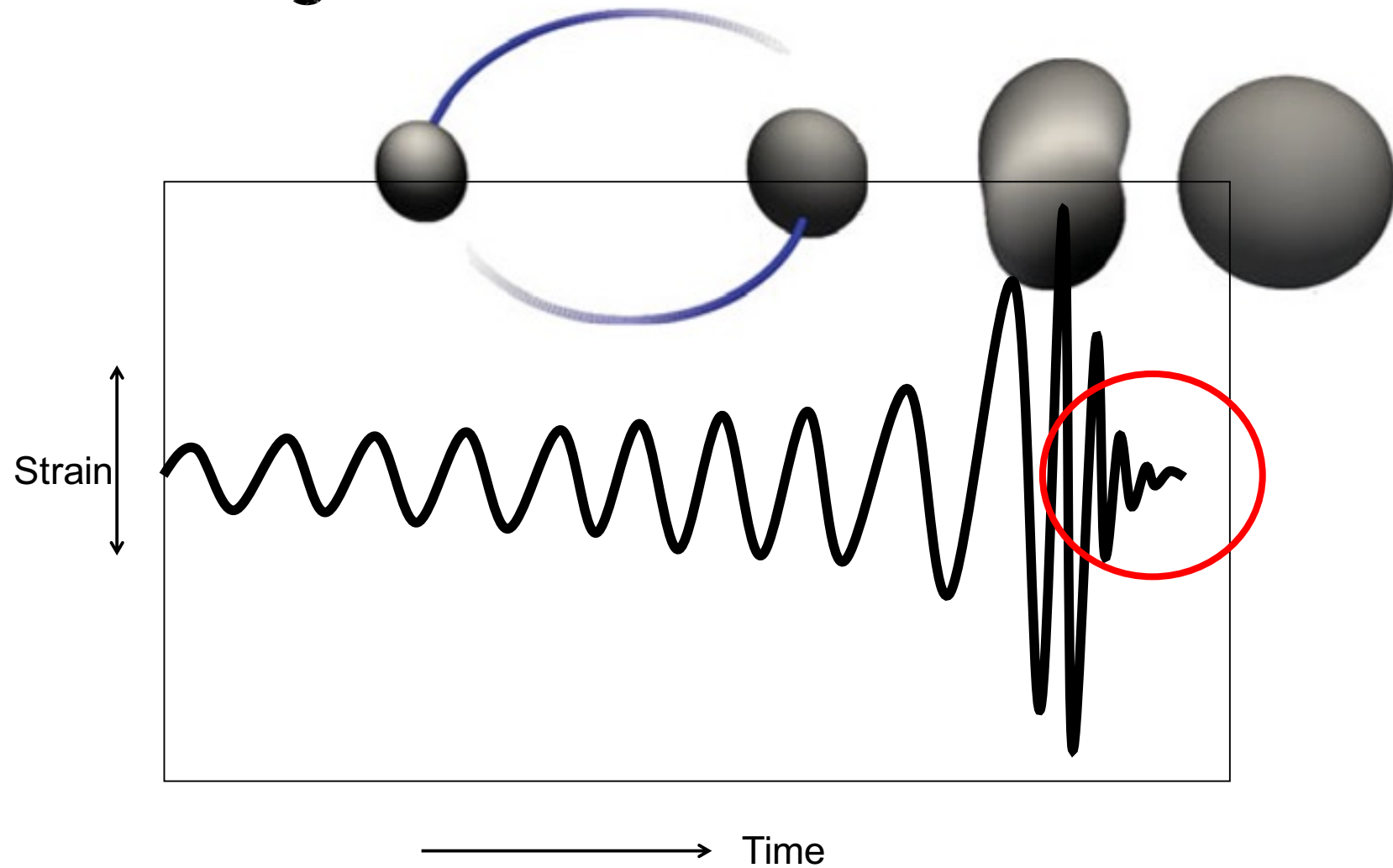
Darkened images

- Δt_{ec} : elapsed coordinate time. We have $\Delta t_{ec}/a_0 \sim e^{a_0^2/nl_p^2}$
- Consider only light rays with $\Delta t_{ec} \leq a_0^3/nl_p^2$ (evaporation timescale)

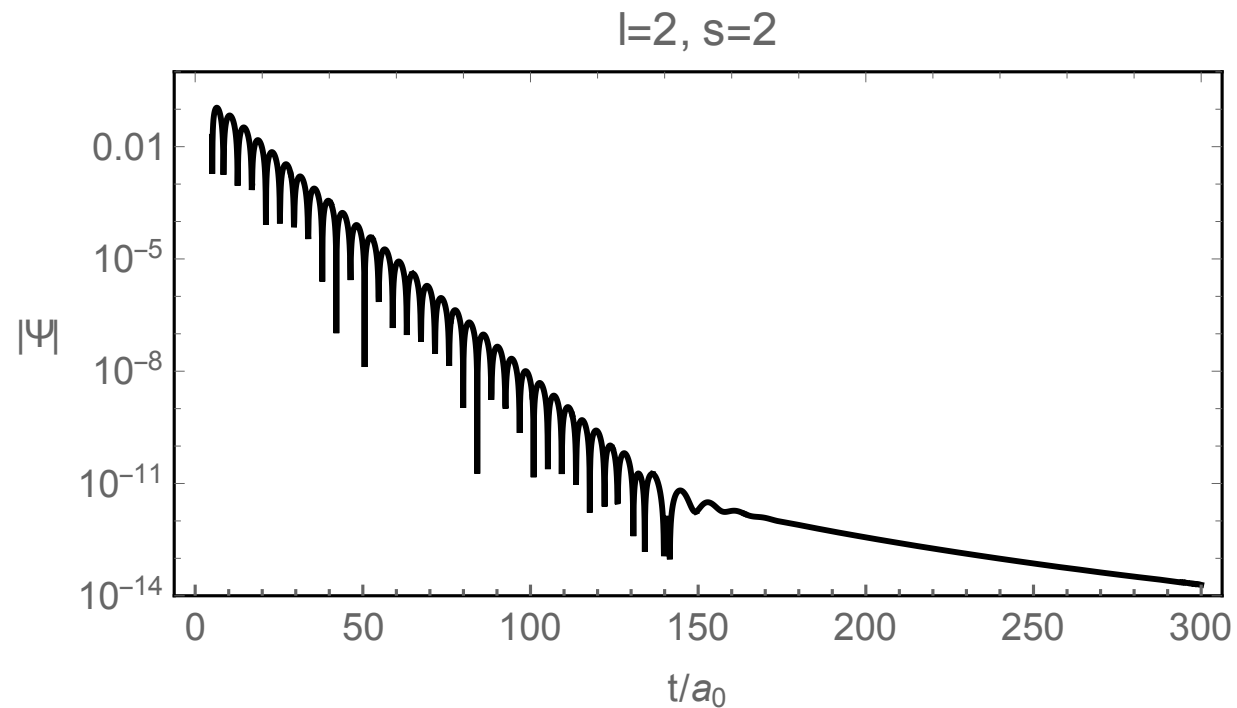
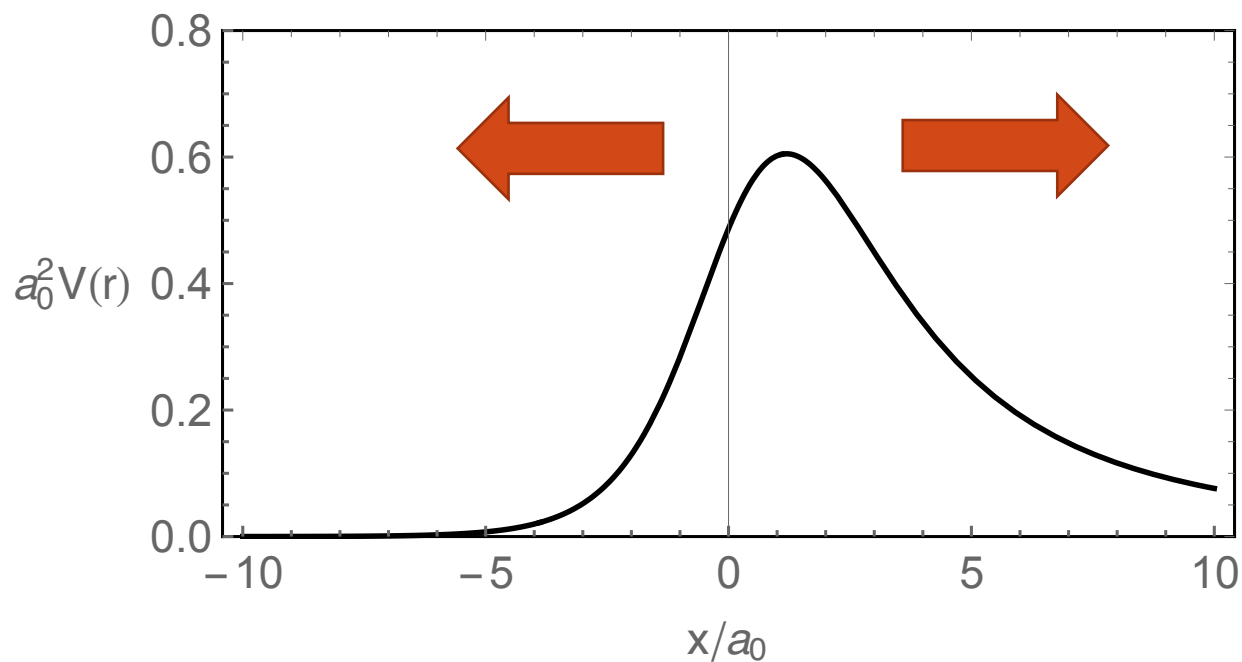


Perfect black hole mimickers!

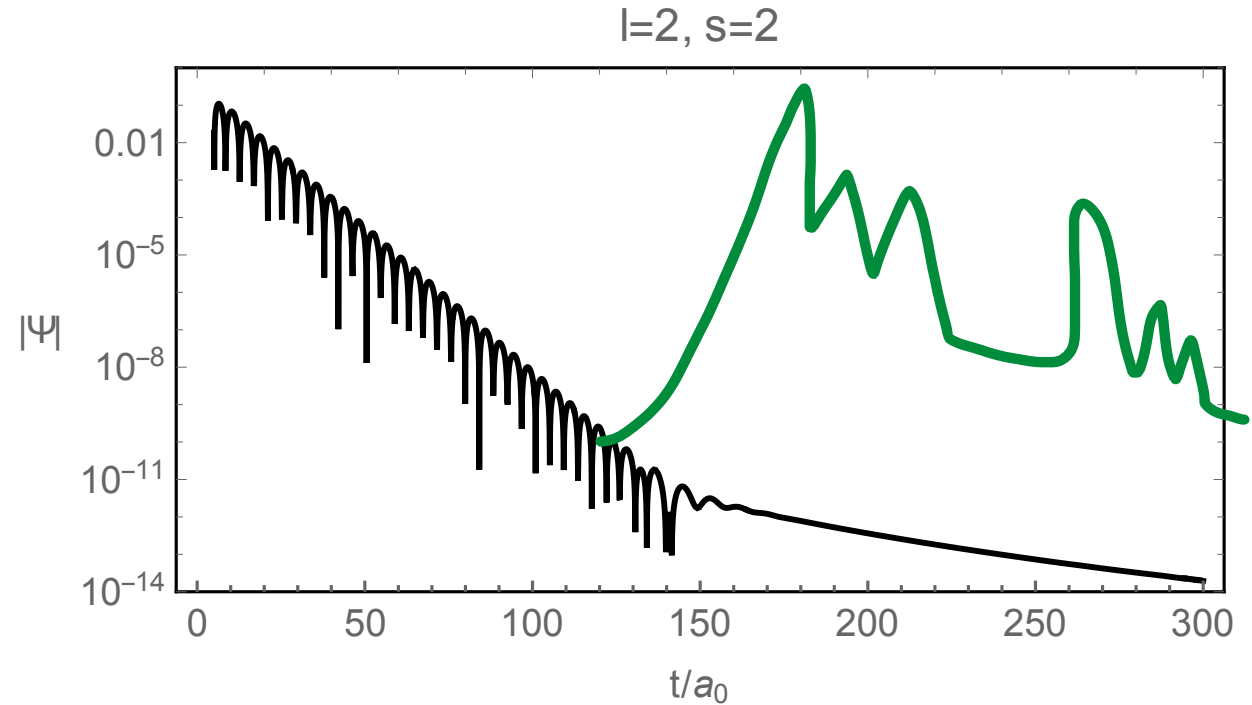
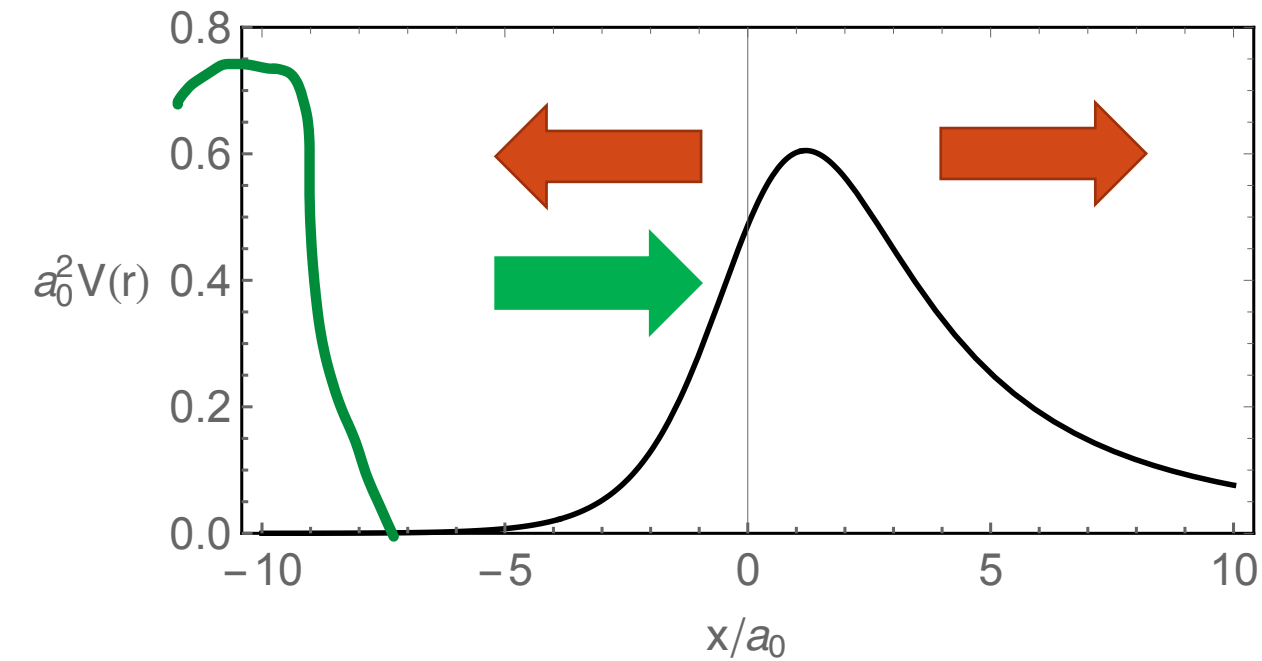
Black hole ringdown



Black hole ringdown

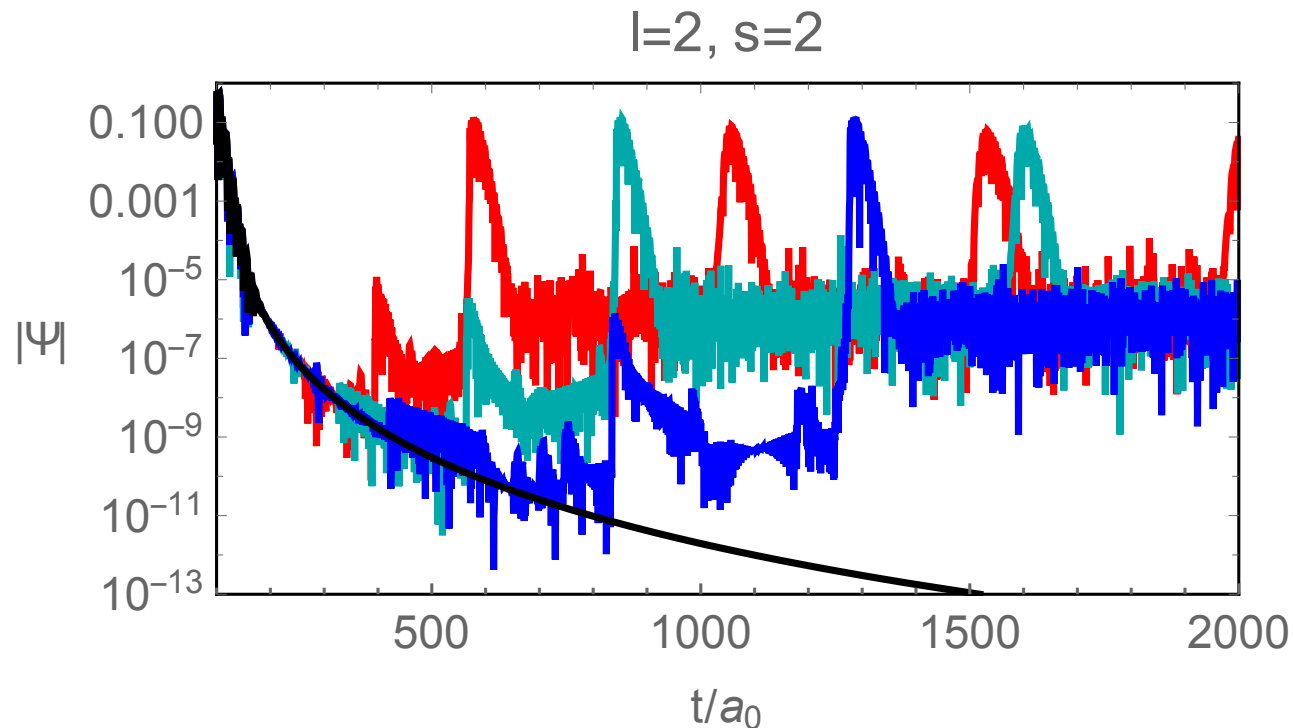


Gravitational wave echoes



QHCO echoes delayed by redshifts

- Series of echoes after typical ringdown
- $\Delta t_{echo}/a_0 \sim e^{a_0^2/nl_p^2}$: time separation between echo peaks



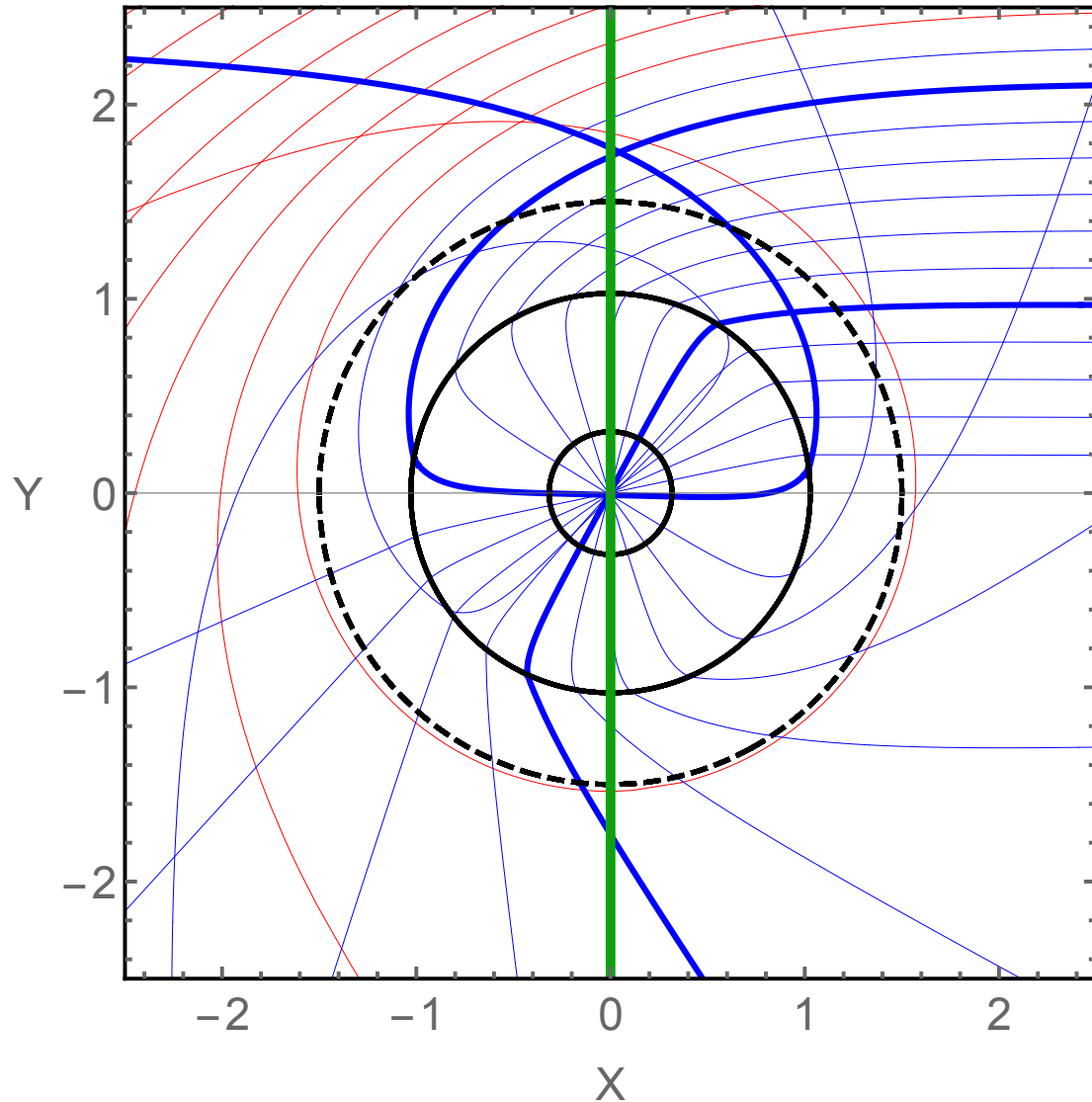
Conclusions

- Bright ring & central darkness do not imply a standard BH
- Darkening effects induced by redshifts are crucial in modeling HCO images
- Redshifts delay GW echoes
- Primordial BHs & dark matter (Shuntaro's talk on early universe)
- QHCO as the end state of evaporation of primordial BHs?
 - e.g. memory burden effects or other UV physics (Pei-Ming's talk)
Dvali (2024) Ho, Kawai, Shao (2024)
- DM candidates?
- Detection through: microlensing, NS capture, interaction with WD, gravitational waves, etc...
Kaiser et al (2025) Carr, Kuhnel (2020)(2022), Green, Kavanagh (2021)

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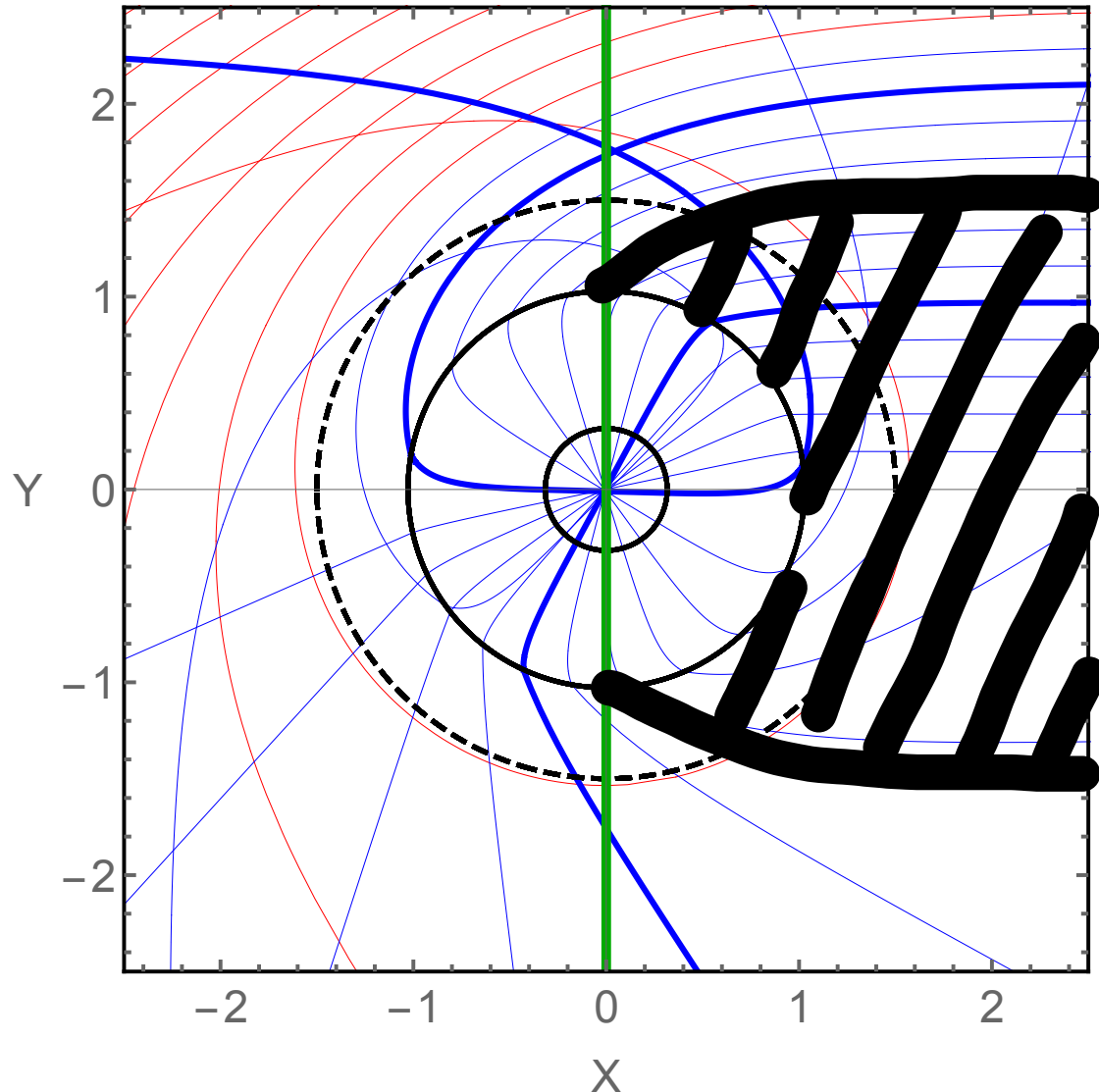
Inner shadow



- Inner shadow boundary:
Direct emission at the horizon



Inner shadow

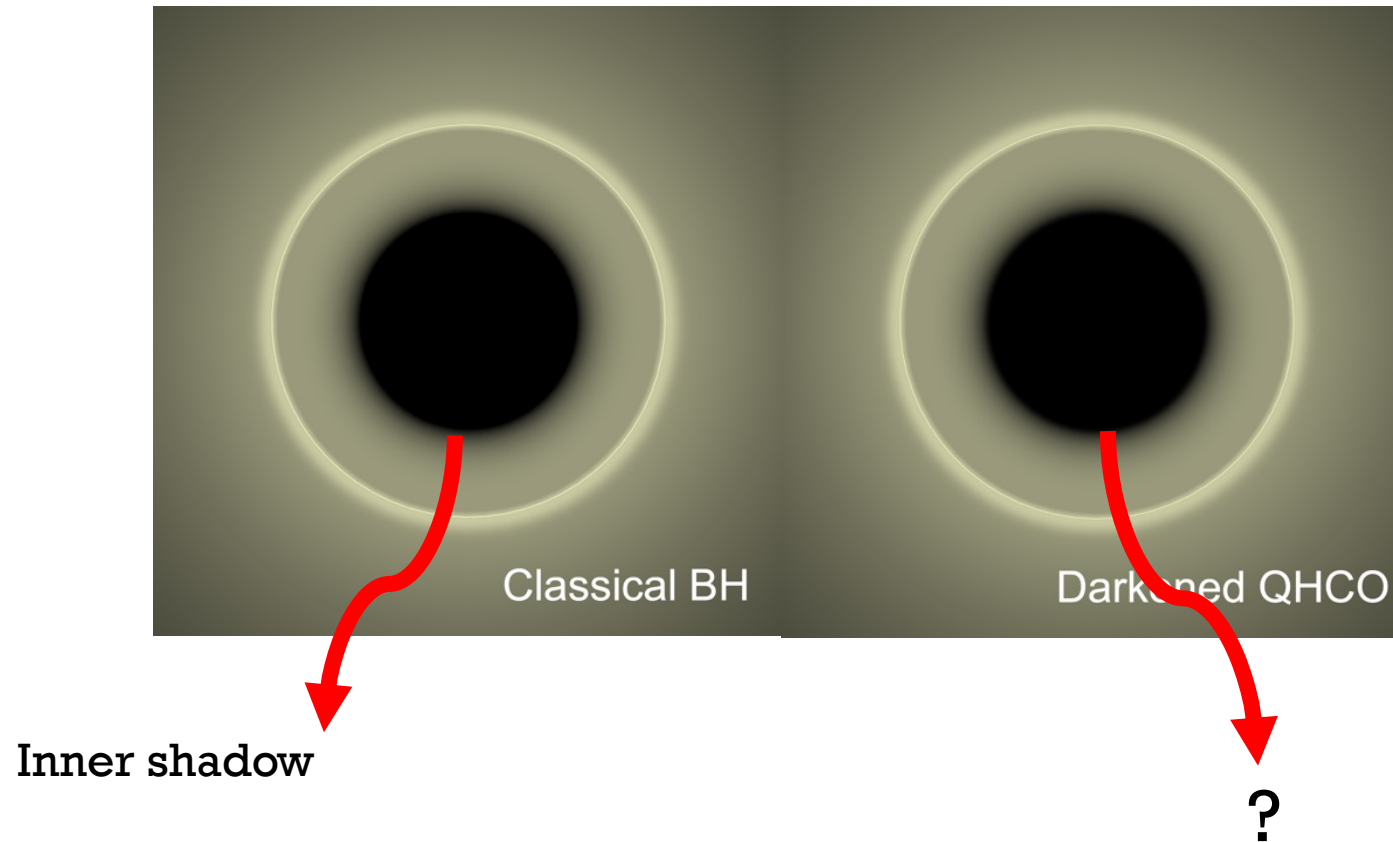


- Inner shadow boundary:
 - Direct emission at the horizon
 - Only valid when emission exists there



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Near the inner shadow: excess intensity

If the light source extends slightly inside the surface R_{out} :

